# $\beta$ matrix models in the multi cut regime

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#### Model definition

## Distributions in $\mathbb{R}^n$ , depending on the function V and $\beta > 0$

$$p_{n,\beta}(\lambda_1,...,\lambda_n) = Z_{n,\beta}^{-1}[V]e^{\beta H(\lambda_1,...,\lambda_n)/2},$$

where H (Hamiltonian) and  $Z_{n,\beta}[V]$  (partition function) are

$$\begin{split} H(\lambda_1,\ldots,\lambda_n) &= -n \sum_{i=1}^n V(\lambda_i) + \sum_{i\neq j} \log |\lambda_i - \lambda_j|, \\ Z_{n,\beta}[V] &= \int e^{\beta H(\lambda_1,\ldots,\lambda_n)/2} d\lambda_1 \ldots d\lambda_n, \quad V(\lambda) > (1+\varepsilon) \log(1+\lambda^2). \end{split}$$

For  $\beta = 1, 2, 4$  it is a joint eigenvalues distribution of real symmetric, hermitian and symplectic matrix models respectively.

#### Marginal densities (correlation functions)

$$p_l^{(n)}(\lambda_1,...,\lambda_l) = \int_{\mathbb{R}^{n-l}} p_{n,\beta}(\lambda_1,...\lambda_l,\lambda_{l+1},...,\lambda_n) d\lambda_{l+1}...d\lambda_n$$

## Linear eigenvalue statistics and characteristic functional

Linear eigenvalue statistics (LES) of the test function h and NCM

$$\mathcal{N}_n[h] = \sum_{i=1}^n h(\lambda_i), \quad N_n[\Delta] = \sharp \{\lambda_i \in \Delta\}/n$$

The moments of LES can be written in terms of correlation functions. In particular,

$$E\{\mathcal{N}_n[h]\} = n \int h(\lambda) p_1^{(n)}(\lambda) d\lambda$$

and  $\operatorname{Var}_n\{\mathcal{N}_n[h]\}\$  can be expressed in terms of  $p_2^{(n)}(\lambda_1,\lambda_2)$  and  $p_1^{(n)}(\lambda_1)$ .

#### Characteristic functional

$$\begin{split} \widetilde{Z}_{n,\beta}[h] = & E_{\beta,n} \Big\{ e^{\beta(\mathcal{N}_n[h] - E\{\mathcal{N}_n[h]\})/2} \Big\} \\ = & Z_{n,\beta}[V - \frac{1}{n}(h - E\{\mathcal{N}_n[h]\})]/Z_{n,\beta}[V], \end{split}$$

# First step for $\beta$ matrix models

### Theorem [Boutet de Monvel, Pastur, S:95; Johansson:98]

If V is a Hölder function, then

$$\log Z_{n,\beta}[V] = \frac{n^2\beta}{2}\mathcal{E}[V] + O(n\log n),$$

where 
$$\mathcal{E}[V] = -\min_{m \in \mathcal{M}_1} \left\{ L[dm, dm] + \int V(\lambda) m(d\lambda) \right\} = \mathcal{E}_V(m^*),$$

$$L[dm, dm'] = \int \log |\lambda - \mu|^{-1} dm(\lambda) dm'(\mu),$$

 $m^*(d\lambda) = \rho(\lambda)d\lambda$  (called the equilibrium measure) has a compact support  $\sigma := \text{supp } m^*$ .

Moreover, if  $h' \in L_2[\sigma_{\varepsilon}]$ 

$$|n^{-1} E\{\mathcal{N}_n[h]\} - (h,m^*)| \leq C n^{-1/2} \log^{1/2} n ||h'||_2^{1/2} ||h||_2^{1/2}$$

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# Motivation to study $\log Z_{n,\beta}[V]$ : universality for $\beta = 1,4$

#### Result of Widom:99

For polynomial V of degree 2m there is  $(2m-1) \times (2m-1)$  matrix  $T_n$  (it can be constructed directly) such that if  $\log \det T_n > -C$  uniformly in n, then the Dyson universality conjecture is true for  $\beta = 1, 4$ 

- $V = \lambda^4/4 + a\lambda^2/2$  [Stojanovich:02],
- $V = \lambda^{2m}$  [Deift,Gioev:07,07a],
- V-real analytic with one interval equilibrium density [S:09,09a].

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## Observation of Stojanovich [St:02]

$$\det(T_n) = \left(\frac{Z_{n,1}[V]Z_{n/2,4}[V]}{Z_{n,2}[V](n/2)!2^n}\right)^2$$

Hence to control  $\det(T_n)$ , it suffices to control  $\log(Z_{n,\beta}/n!)$  for  $\beta = 1, 2, 4$  up to O(1) terms.

## Results for one cut potentials

#### Theorem [Johansson:98] CLT for LES in the one cut case

V is polynomial,  $\sigma = [a, b]$ , and  $\rho$  is "generic". Then for any  $h : \mathbb{R} \to \mathbb{R}$  with  $||h^{(4)}||_{\infty}, ||h'||_{\infty} \leq \log n$ 

$$\widetilde{Z}_{n,\beta}[h] = \exp\Big\{\frac{\beta}{2}\Big(\big(\frac{2}{\beta}-1\big)(h,\nu) + \frac{1}{4}(\overline{D}_{\sigma}h,h)\Big)\Big\}\Big(1 + n^{-1}O\big(||h^{(4)}||_{\infty}^{3}\big)\Big)$$

where the "variance operator"  $\overline{\mathrm{D}}_{\sigma}$  and the measure  $\nu$  have the form

$$(\overline{D}_{\sigma}h, h) = \int_{\sigma} \frac{h(\lambda)d\lambda}{\pi^{2}X^{1/2}(\lambda)} \int_{\sigma} \frac{h'(\mu)X^{1/2}(\mu)d\mu}{\lambda - \mu}, X_{\sigma}(\lambda) = (b - \lambda)(\lambda - a)$$
$$(\nu, h) := \frac{1}{4}(h(b) + h(a)) - \frac{1}{2\pi} \int_{\sigma} \frac{h(\lambda)d\lambda}{X^{1/2}(\lambda)} + \frac{1}{2}(D_{\sigma}\log P, h)$$

#### Remark

 $D_{\sigma}$  is "almost"  $\mathcal{L}_{\sigma}^{-1}$ , where  $\mathcal{L}_{\sigma}$  is the integral operator defined by the kernel  $\log |\lambda - \mu|^{-1}$  for the interval  $\sigma$ 

#### Theorem [Kriecherbauer, S:10]

• For h = 0

$$\begin{split} \log(Z_{n,\beta}/n!) = & \frac{\beta n^2}{2} \mathcal{E}[V] + F_{\beta}(n) + n \left(\frac{\beta}{2} - 1\right) \left((\log \rho, \rho) - 1 - \log 2\pi\right) \\ & + r_{\beta}[\rho] + O(n^{-1}), \end{split}$$

where  $F_{\beta}(n)$  corresponds to the linear, logarithmic and zero order terms of the expansion in n of  $\log Z_{n,\beta}[V^*]$  for  $V^*(\lambda) = \lambda^2/2$ :

$$F_{\beta}(n) = n\left(\frac{\beta}{2} - 1\right)\left(\log\frac{n\beta}{2} - \frac{1}{2}\right) + n\log\frac{\sqrt{2\pi}}{\Gamma(\beta/2)} - c_{\beta}\log n + c_{\beta}^{(1)},$$

where  $c_{\beta} = \frac{\beta}{24} - \frac{1}{4} + \frac{1}{6\beta}$  and  $c_{\beta}^{(1)}$  is some depending only on  $\beta$  constant (for  $\beta = 2$ ,  $c_{\beta}^{(1)} = \zeta'(1)$ )

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## Other results for one cut potentials:

- [Albeverio, Pastur, S:01] expansion of the first and the second correlators for one-cut real analytic V and  $\beta = 2$ ;
- [Ercolani, McLaughlin:03] expansion of  $\log Z_{n,\beta}[V]$  for polynomial one-cut V and  $\beta = 2$ ;
- **3** [Borot, Guionnet:11] expansion of all correlators and  $\log Z_{n,\beta}[V]$  for one-cut real analytic V and any  $\beta$ .

## CLT and expansions for multi - cut case. Results.

- [Chekhov, Eynard: 06, Eynard: 09] formal expansions for multi-cut V and any  $\beta$ ;
- ② [Pastur:07] derivation of CLT from OP-asymptotics of [Deift at al:99] for real analytic h and  $\beta = 2$ ;

# Idea from the mean field theory of statistical mechanics

Consider the Hamiltonian

$$H_n(\bar{\sigma}) = H_n^*(\bar{\lambda}) + \frac{1}{2} \Big( \sum_j \varphi(\lambda_j) \Big)^2,$$

where  $H_n^*$  is the Hamiltonian for which we are able to find the  $\log Z_n^*(u)$  up to the order  $O(n^{-k})$ 

$$f_n^*(u) = \beta^{-1} \log Z_n^*(u) = \beta^{-1} (\log Z_n^*(0) + d\frac{u^2}{2} + \sum d_k(u)n^{-k})$$

where

$$Z_n^*[u] := \int d\bar{\lambda} e^{\beta(H_n(\bar{\lambda}) + nv \sum \varphi(\lambda_i) + u \sum (\varphi(\lambda_i) - \langle \varphi \rangle))}.$$

We would like to find the partition function of  $H_n$ :

$$Z_n = \int d\bar{\lambda} e^{\beta H_n(\bar{\lambda})}$$

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• Introduce the "approximate" Hamiltonian

$$H_n^{(a)}(\bar{\lambda}) = H_n^*(\bar{\lambda}) + nv \sum_j \varphi(\lambda_i) - n^2 \frac{v^2}{2}, \quad v = \langle \varphi \rangle_{H_n}$$

$$H_{\mathrm{n}} = H_{\mathrm{n}}^{\mathrm{(a)}}(\bar{\lambda}) + \frac{1}{2} \sum_{\mathrm{i},\mathrm{i}} (\varphi(\lambda_{\mathrm{i}}) - \langle \varphi \rangle_{\mathrm{H}_{\mathrm{n}}}) (\varphi(\lambda_{\mathrm{j}}) - \langle \varphi \rangle_{\mathrm{H}_{\mathrm{n}}})$$

• Use the Hubbard-Stratonovich transformation

$$Z_n = \sqrt{\frac{\beta}{2\pi}} \int d\bar{\lambda} \int du e^{\beta H_n^{(a)}(\bar{\lambda}) + u\beta \sum_j (\phi(\lambda_i) - \langle \phi \rangle_{H_n}) - \beta u^2/2}$$

• Take the integral with respect to  $\bar{\lambda}$  first. We obtain

$$Z_n = \int du e^{\beta f_n^*(u) - \beta u^2/2}$$

• Expand  $f_n^*(u)$  in the series with respect to  $n^{-1}$  and live at the exponent only O(1) quadratic term. Taking the integrals with respect to u we obtain the expansion for  $Z_n$ .

Assumptions and the restriction of the integration domain

#### Assumptions

• V is real analytic,

$$\sigma = \bigcup_{\alpha=1}^{q} \sigma_{\alpha}, \quad \mu_{\alpha} = \int_{\sigma_{\alpha}} \rho_{\alpha}(\lambda) d\lambda, \quad \rho_{\alpha} := 1_{\sigma_{\alpha}} \rho.$$

2 V is of generic behavior.

Replace the integration domain in the definition of  $Z_{n,\beta}[V]$  from  $\mathbb{R}$  to  $\sigma_{\varepsilon}$ , where

$$\sigma_{\varepsilon} = \bigcup_{\alpha=1}^{q} \sigma_{\alpha,\varepsilon}, \quad \sigma_{\alpha,\varepsilon} \cap \sigma_{\alpha+1,\varepsilon} = \emptyset$$

Then, according to the result of [Pastur,S:07],  $Z_{n,\beta}[V]$  will be changed by  $(1 + O(e^{-nc}))$  factor.

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## Construction of the "approximate" Hamiltonian

Below we will use the notation

$$\chi_{\alpha}(\lambda) = 1_{\sigma_{\alpha,\varepsilon}}(\lambda),$$

Then for our domain  $1 = \sum_{\alpha} \chi_{\alpha}(\lambda)$  and we can write  $H(\lambda)$  as

$$\begin{split} H(\overline{\lambda}) &= -n \sum_{i=1}^{n} V(\lambda_{i}) + \sum_{i \neq j, \alpha, \alpha' = 1}^{q} \chi_{\alpha}(\lambda_{i}) \chi_{\alpha'}(\lambda_{j}) \log |\lambda_{i} - \lambda_{j}| \\ &= -n \sum_{i=1}^{n} V(\lambda_{i}) + \sum_{i \neq j} \sum_{\alpha = 1}^{q} \chi_{\alpha}(\lambda_{i}) \chi_{\alpha}(\lambda_{j}) \log |\lambda_{i} - \lambda_{j}| \\ &+ \sum_{\substack{i,j=1 \\ \alpha \neq \alpha'}}^{n} \int \log |\lambda - \mu| \chi_{\alpha}(\lambda) \chi_{\alpha'}(\mu) \delta_{\lambda_{i}}(\lambda) \delta_{\lambda_{j}}(\mu) d\lambda d\mu = H^{*} \end{split}$$

Then write under the integral sign

$$\delta_{\lambda_i}(\lambda) = \delta_{\lambda_i}(\lambda) - \langle \delta_{\lambda_i}(\lambda) \rangle + \langle \delta_{\lambda_i}(\lambda) \rangle = \Delta_i(\lambda) + \langle \underline{\delta_{\lambda_i}(\lambda)} \rangle$$

# Construction of the "approximate" Hamiltonian

$$\begin{split} H(\overline{\lambda}) = & H^*(\overline{\lambda}) + 2n \sum_{j=1}^n V_{\alpha}^{(a)}(\lambda_i) - n^2 \Sigma^* \\ & + \sum_{\substack{i,j=1\\ \alpha \neq \alpha'}}^n \int \log|\lambda - \mu| \chi_{\alpha}(\lambda) \chi_{\alpha'}(\mu) \Delta_i(\lambda) \Delta_j(\mu) \mathrm{d}\lambda \mathrm{d}\mu \\ = & H_a(\overline{\lambda}) + \Delta H(\overline{\lambda}), \end{split}$$

where, taking into account that  $\langle \delta_{\lambda_i}(\lambda) \rangle = \rho(\lambda)$ , we obtain that the "effective potentials"  $V_{\alpha}^{(a)}$  and  $\Sigma^*$  have the form

$$V_{\alpha}^{(a)}(\lambda) = \chi_{\alpha}(\lambda) \sum_{\alpha' \neq \alpha} \int \log|\lambda - \mu| \chi_{\alpha'}(\mu) \rho(\mu) d\mu$$
$$\Sigma^* := \sum_{\alpha \neq \alpha'} \int_{\sigma_{\alpha}} d\lambda \int_{\sigma_{\alpha'}} d\mu \log|\lambda - \mu| \rho(\lambda) \rho(\mu)$$

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It is easy to see that

$$Z_{n,\beta}[V]/n! = \sum_{|\overline{n}|=n} \frac{\int 1_{\overline{n}}(\overline{\lambda}) e^{\beta H(\overline{\lambda})/2}}{n_1! \dots n_q!} = \sum_{|\overline{n}|=n} \frac{\int 1_{\overline{n}}(\overline{\lambda}) e^{\beta (H_a(\overline{\lambda}) + \Delta H(\overline{\lambda}))/2}}{n_1! \dots n_q!}$$

where  $\overline{n} := (n_1, \dots, n_q), |\overline{n}| := \sum_{\alpha=1}^q n_\alpha$ , and  $1_{\overline{n}}(\overline{\lambda})$  is the indicator of the configurations of  $\lambda_1, \dots, \lambda_n$ , such that

 $\begin{array}{l} \lambda_1,\ldots,\lambda_{n_1}\in\sigma_{1,\varepsilon},\,\lambda_{n_1+1},\ldots,\lambda_{n_1+n_2}\in\sigma_{2,\varepsilon},\ldots\,\,\lambda_{n-n_q+1},\ldots,\lambda_q\in\sigma_{q,\varepsilon}\\ \mathrm{Choose}\,\,\mathrm{M}=[\log^2 n]\,\,\mathrm{and}\,\,\mathrm{represent}\,\,\Delta\mathrm{H}(\overline{\lambda})\mathbf{1}_{\overline{n}}(\overline{\lambda})\,\,\mathrm{as} \end{array}$ 

$$\begin{split} &\mathbf{1}_{\overline{n}}(\overline{\lambda}) \sum_{\substack{j,j'=1\\\alpha \neq \alpha'}}^{n} \sum_{k,m=1}^{M} L_{k,m}^{(\alpha,\alpha')} \Big( p_k^{(\alpha)}(\lambda_j) - \frac{n}{n_\alpha} c_k^{(\alpha)} \Big) \Big( p_m^{(\alpha')}(\lambda_{j'}) - \frac{n}{n_{\alpha'}} c_k^{(\alpha)} \Big) \\ &+ O(e^{-c \log^2 n}), \quad c_k^{(\alpha)} := (p_k^{(\alpha)}, \rho \mathbf{1}_{\sigma_\alpha}) \end{split}$$

where  $L_{k,m}^{(\alpha,\alpha')}$  are the Fourier coefficient of the function

$$\log |\lambda - \mu| \chi_{\alpha}(\lambda) \chi_{\alpha'}(\mu)$$
 with respect to the basis  $\{p_k^{(\alpha)}(\lambda) p_m^{(\alpha')}(\mu)\}$ 

## Main steps of the proof

- For each  $\bar{\mathbf{n}}$  we linearize  $\Delta \mathbf{H}(\overline{\lambda})\mathbf{1}_{\overline{\mathbf{n}}}(\overline{\lambda})$ , using the Hubbard-Stratonovich transformation. This adds to  $\mathbf{H}^{(a)}(\overline{\lambda})\mathbf{1}_{\overline{\mathbf{n}}}(\overline{\lambda})$  the additional potential  $\mathbf{h}_{\alpha}[\bar{\mathbf{u}}]$ , depending on the integration parameters  $\bar{\mathbf{u}}$ . Then for each  $\sigma_{\alpha}$  we are in the situation of Theorem 3.
- Apply Theorems 2,3 to find  $Z_{n_{\alpha},\beta}[V+2V_{\alpha}^{(a)}+\frac{1}{n}h_{\alpha}[\bar{u}]]$ . We obtain the quadratic form of  $\bar{u}$  in the exponent.
- Integrate with respect to  $\bar{\mathbf{u}}$ . We obtain the expansion for the initial partition function.

## Important definitions

$$X_{\sigma}(\lambda) = \prod_{\alpha=1}^{q} (b_{\alpha} - \lambda)(\lambda - a_{\alpha})$$

#### Definition of Q

$$Q = \{Q_{\alpha\alpha'}\}_{\alpha,\alpha'=1}^{q}, \quad Q_{\alpha\alpha'} = (\mathcal{L}\psi^{(\alpha)}, \psi^{(\alpha')}),$$

where  $\psi^{(\alpha)}(\lambda) = p_{\alpha}(\lambda)X^{-1/2}(\lambda)1_{\sigma}$  ( $p_{\alpha}$  is a polynomial of degree q-1) is a unique solution of the system of equations

$$(\mathcal{L}\psi^{(\alpha)})_{\alpha'} = \delta_{\alpha\alpha'}, \quad \alpha' = 1, \dots, q.$$

(harmonic measure of  $\sigma_{\alpha}$  with respect to  $\mathbb{C} \setminus \sigma$ )

## Definition of I[h]

$$I[h] = (I_1[h], \dots, I_q[h]), \quad I_{\alpha}[h] := \sum_{\alpha'} \mathcal{Q}_{\alpha\alpha'}^{-1}(h, \psi^{(\alpha')}).$$

#### Main results

#### Theorem 3 [S:12]

Let the potential V satisfy conditions C1-C2. Then

$$\begin{split} \widetilde{Z}_{n,\beta}[h] = & e^{\frac{\beta}{8}(\mathcal{D}h,h) + \left(\frac{\beta}{2} - 1\right)(\mathcal{G}\nu,h)} \frac{\Theta(\overline{I}[h];\{n\overline{\mu}\})}{\Theta(0;\{n\overline{\mu}\})} \\ & \times \left(1 + O\left(n^{-\kappa}(||h'||_{\infty}||h^{(6)}||_{\infty}^{2})\right)\right), \end{split}$$

where the operators  $\mathcal{D}$ ,  $\mathcal{G}$ ,  $\widetilde{\mathcal{L}}$  are defined in terms of  $\mathcal{L}$  and  $\oplus D_{\sigma_{\alpha}}$ 

$$\begin{split} \Theta(I[h]; \{n\bar{\mu}\}) := \sum_{n_1 + \dots + n_q = n_0} \exp\Big\{ -\frac{\beta}{2} \Big( \mathcal{Q}^{-1} \Delta \bar{n}, \Delta \bar{n} \Big) + \frac{\beta}{2} (\Delta \bar{n}, I[h]) \\ + \Big( \frac{\beta}{2} - 1 \Big) (\Delta \bar{n}, I[\log \overline{\rho}]) \Big\}, \end{split}$$

$$\{n\bar{\mu}\} = (\{n\mu_1\}, \dots, \{n\mu_q\}), \ (\Delta\bar{n})_{\alpha} = n_{\alpha} - \{n\mu_{\alpha}\}, \ n_0 = \sum^q \{n\mu_{\alpha}\},$$

 $\bullet$  for h = 0 we have

$$\begin{split} Z_{n,\beta}[V] = & \mathcal{S}_{n,\beta}[V] \frac{\exp\left\{\frac{2}{\beta}\left(\frac{\beta}{2}-1\right)^2(\widetilde{\mathcal{L}}\mathcal{G}\nu,\nu)\right\}}{\det^{1/2}(1-\overline{D}\widetilde{\mathcal{L}})} \Theta(0;\{n\bar{\mu}\})(1+O(n^{-\kappa})), \\ \mathcal{S}_{n,\beta}[V] = & \exp\left\{\frac{n^2\beta}{2}\mathcal{E}[V] + F_{\beta}(n) + n(\frac{\beta}{2}-1)\big((\log\rho,\rho) - 1 - \log 2\pi\big) - c_{\beta}(q-1)\log n + \sum_{\alpha=1}^{q}(r_{\beta}[\mu_{\alpha}^{-1}\rho_{\alpha}] - c_{\beta}\log\mu_{\alpha})\right\}, \end{split}$$

#### Theorem 4 [S:in prep]

Under the conditions C1-C2  $Q_{n,\beta}[V]$  admits the asymptotic expansion in  $n^{-j}$  with quasi-periodic in n coefficients  $q_{\beta,j}[n]$ :

$$Z_{n,\beta}[V] = \mathcal{S}_{n,\beta}[V] \frac{\exp\left\{\frac{2}{\beta}\left(\frac{\beta}{2}-1\right)^2\left(\widetilde{\mathcal{L}}\mathcal{G}\nu,\nu\right)\right\}}{\det^{1/2}(1-\overline{D}\widetilde{\mathcal{L}})} \Theta(0;\{n\bar{\mu}\}) \sum_{j=1} n^{-j}q_{\beta,j}[n],$$

#### Corollaries

#### Corollary 1

Theorem 3 yields that the fluctuations of  $\mathcal{N}_n[h]$  for generic h are non Gaussian. They are Gaussian, if there exists some c such that

$$I_{\alpha}[h] = c, \quad \alpha = 1, \dots, q; \quad \Leftrightarrow \quad (h - c, \psi^{(\alpha)}) = 0, \quad \alpha = 1, \dots, q.$$

#### Remark 1

The operator  $\mathcal{D}$ , which appears in the place of the "variance operator" in the multi-cut case, differs from  $\mathcal{L}^{-1}$  only by the finite rank perturbation. This perturbation provides, in particular, that  $\mathcal{D}f = 0$ , if  $f(\lambda) = \text{const}$ ,  $\lambda \in \sigma$ 

$$(\mathcal{D}h, h) = \frac{1}{\pi^2} \int_{\sigma} \frac{h(\lambda) d\lambda}{X^{-1/2}(\lambda)} \int_{\sigma} \frac{h'(\mu) X^{1/2}(\mu) d\mu}{(\lambda - \mu)} - ?$$

# Corollaries of Theorem 3 for the moments of $\mathcal{N}_n[h]$

## Expectation of $\mathcal{N}_n[h]$

For any  $\beta$  in the multi cut case we obtain  $O(n^{-1})$  correction to  $E_{\beta,n}\{n^{-1}\mathcal{N}_n[h]\}$ :

$$E_{\beta,n}\{n^{-1}\mathcal{N}_n[h]\} - (h,\rho) = \frac{1}{n} \left[ \left( \frac{\beta}{2} - 1 \right) (\mathcal{G}\nu,h) + \sum_{\alpha=1}^q I_{\alpha}[h] c_{\alpha}(n) \right] + O\left( \frac{1}{n^{1+\delta}} \right),$$

## Variance of $\mathcal{N}_{n}[h]$

$$\operatorname{Var}\{\mathcal{N}_{\mathbf{n}}[\varphi]\} = \frac{\beta}{8}(\mathcal{D}\mathbf{h}, \mathbf{h}) + \sum_{\alpha=1}^{q} d_{\alpha, \alpha'}(\mathbf{n}) I_{\alpha}[\mathbf{h}] I_{\alpha'}[\mathbf{h}]$$

where  $c_{\alpha}(n)$  and  $d_{\alpha,\alpha'}(n)$  are quasi-periodic functions (derivatives of  $\Theta$ -function above).

# Corollaries from Theorem 3 for the universality of local regimes

- bulk universality for  $\beta = 1, 4$ ;
- 2 edge universality for  $\beta = 1, 4$ ;
- ullet bulk universality for any eta could be reduced to the universality for the one-cut case for  $V + n^{-1}h$  with  $||h'||_{\infty} \leq \log n$ .